

Single-Longitudinal-Mode Fiber Ring Lasers With Taper-Coupled Double-Microsphere-Cavities

Hongdan Wan¹, Linqian Liu, Zuoqin Ding, Jie Wang, Kaidi Lu, Zuxing Zhang¹, and Lin Zhang

Abstract—This letter proposes and demonstrates a fiber ring laser using taper-coupled double-microsphere-cavities (DMCs) to achieve single-longitudinal-mode operation. Whispering-gallery-mode (WGM) intensity distributions and transmission spectra of the DMC with different microsphere diameters are investigated both theoretically and experimentally. Due to the Vernier effect, the DMC can produce WGM spectra with a higher extinction ratio, a higher side-mode-suppression ratio, a larger free spectral range, and a narrower bandwidth, as compared with a single-microsphere cavity. A single-longitudinal-mode fiber ring laser operating near $1.5 \mu\text{m}$ with a bandwidth of $< 0.01 \text{ nm}$ and a signal-to-background ratio of about 60 dB is demonstrated using the taper-coupled DMC as an all-fiber mode selector.

Index Terms—Fiber laser, single-longitudinal-mode, taper-coupled double-microsphere-cavities, whispering-gallery-mode, Vernier effect.

I. INTRODUCTION

AS APPLIED in optical sensing and communication systems, single-longitudinal-mode (SLM) narrow-linewidth fiber lasers are of great interest to be developed because of their high coherence, flexible wavelengths and compact structures [1]–[3]. Various narrow-bandwidth mode-selection methods have been developed and introduced into fiber laser cavities in order to achieve SLM operation, such as fiber gratings [4], [5], saturated absorbers [6], tandem all-Fiber Fabry-Pérot micro-cavities and gain-controlled active compound cavities, which we previously proposed in [7] and [8]. However, most of these filters have BWs boarder than 1 GHz [4]–[7]. Recently, an ultra-narrow-bandwidth mode-selection method based on a high-quality (Q) whispering-gallery-mode (WGM) CaF_2 microsphere cavity

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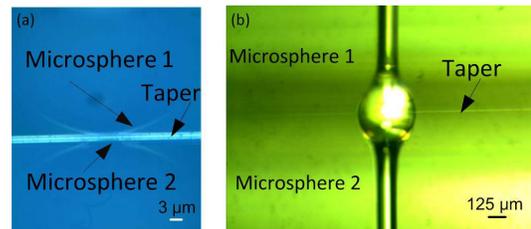


Fig. 1. Microscope images of the taper-coupled DMC: (a) Observed from the horizontal orientation; (b) Observed from the vertical orientation.

was proposed [9]. A SLM laser output with a 650 Hz linewidth was obtained at 1550 nm. However, manufacturing such a CaF_2 microsphere out of a crystal is time- and labor-consuming. Compared to the CaF_2 microsphere, a SiO_2 microsphere can be easily fabricated by utilizing the surface tension effect [10], [11]. In [11], we proposed and demonstrated a SLM fiber ring laser using a narrow WGM spectrum produced by a single SiO_2 microsphere, of which the Q factor of the microsphere highly depends on the single-microsphere-cavity's (SMC) coupling loss and surface quality.

In this letter we propose a SLM selection method based on narrow-bandwidth WGM spectra produced by taper-coupled double-microsphere-cavities (DMC), which consists of two SiO_2 microspheres coupling in parallel with a fiber taper. Based on total internal reflection, the WGMs are confined inside the DMC, where the Vernier effect results in WGM spectra with a higher extinction ratio (ER), an increased side-mode-suppression-ratio (SMSR), a compressed BW and a larger free spectral range (FSR), as compared to the SMC. A $1.5 \mu\text{m}$ -wavelength SLM fiber ring laser with a 3dB bandwidth of 0.01 nm is experimentally demonstrated using the taper -coupled DMC as the all-fiber mode-selector.

II. EXPERIMENTAL SETUPS

Figure 1 shows the microscope image of the taper-coupled DMC with microsphere diameters D_1 of about $198.4 \mu\text{m}$ and D_2 of about $198.3 \mu\text{m}$, observed by two high-resolution microscope systems in the horizontal and vertical directions. The microsphere diameters are measured at the equatorial plane by the microscope system consisting of a 500-megapixel charge-coupled-device (CCD) camera. The highest resolution and the largest total magnification of the measurement are about $0.6 \mu\text{m}$ and 1600, respectively. The two microspheres are fabricated by melting the end of a single-mode optical fiber using a CO_2 laser. They are placed on two opposite

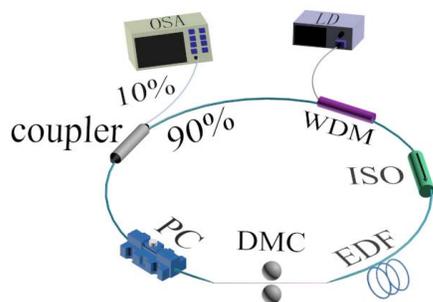


Fig. 2. Experimental setup of the proposed SLM fiber laser stabilized by the taper-coupled DMC.

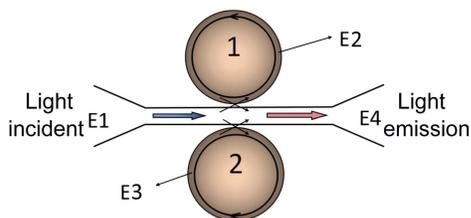


Fig. 3. Light propagation schematic in the taper-coupled DMC.

sides of the tapered fiber. The gap spacing between the microsphere and the taper is controlled by electromechanical three-dimensional (3D) X-Y-Z translation stages with a 20 nm resolution. WGMs are excited by a wavelength-tunable laser (Agilent 81960A, wavelength near 1550 nm) through the taper and detected by a photodiode detector.

Figure 2 shows the experimental setup of the proposed SLM, a narrow-bandwidth fiber ring laser stabilized by the taper-coupled DMC. The fiber ring cavity consists of a 980 nm pump laser diode (LD), a 980/1550 nm wavelength-division multiplexer (WDM), an isolator (ISO), an erbium-doped fiber (EDF) with a length of 1.5 m, a 90:10 fiber coupler, a polarization controller (PC) and the taper-coupled DMC. The pump LD is injected into the EDF through the 980 nm port of the WDM. The two microspheres are placed in parallel on opposite sides of the fiber taper, acting as a SLM narrow-bandwidth mode-selector in the fiber ring cavity. The optical coupler is used to offer feedback and laser output for detection: with a 90% feedback into the ring cavity and a 10% output, measured by an optical spectrum analyzer (OSA), with a resolution of 0.02 nm (YOKOGAWA, AQ6370D).

III. THEORETICAL ANALYSIS

A. Light Propagation Schematic in the Taper-Coupled DMC

Figure 3 shows the schematic of the light propagation inside the taper-coupled DMC, in which, E1 is the incident light intensity, E2 is the WGM light intensity in microsphere 1 (above), E3 is the WGM light intensity in microsphere 2 (below), E4 is the output light intensity. Light is coupled into microsphere 1 via the taper, then coupled into microsphere 2 via the same fiber taper, and vice versa. The WGM resonant effect is enhanced inside the two microspheres through paralleled coupling via the same fiber taper and results

in WGM spectra with a narrowed 3dB bandwidth (BW). The theoretical analysis is given below.

B. The Vernier Effect in the Taper-Coupled DMC

The Vernier effect has been explored in various active microcavities [12]–[14]. For the SMC, its FSR is given as follows:

$$F = \lambda^2 c / (n_{eff} \cdot L) \quad (1)$$

where n_{eff} is the effective refractive index, λ_c is the resonant wavelength, L is the cavity round-trip length. Different microsphere cavities have different FSR values due to their different diameters ($D's$). The FSR of the DMC is given as [15]:

$$F = N F_1 = (N + 1) F_2 \quad (2)$$

where F_1 and F_2 are the FSRs of the two cavities, N is a positive integer, then we obtain:

$$F = F_1 F_2 / |F_2 - F_1| \quad (3)$$

From (3) we see that due to the Vernier effect, the taper-coupled DMC has a $F_2 / |F_2 - F_1|$ times larger FSR than the SMC, which is better for side-mode suppression and narrow-bandwidth SLM selection as compared to the SMC.

C. Comparison Between the SMC and the DMC

The mode-field patterns in the microspheres are simulated by Finite-Difference Time-Domain (FDTD) method previously reported in [16]. Here we detail the simulation parameters of the taper-coupled DMC: The diameters of the two microspheres are the same: $D_1 = D_2 = 35 \mu\text{m}$. The diameters chosen in the simulation are significantly smaller than those in the experiment. Thus, for phase-matched coupling, the taper diameter simulated is different ($1.4 \mu\text{m}$) from that employed in the experiment, which is about $3 \sim 4 \mu\text{m}$. The gap spacing between the microsphere and the taper is set to $0.1 \mu\text{m}$ to achieve phase-matched coupling. The material refractive indices of the taper and the microsphere are 1.46. The light source is transverse-electric (TE) polarized. Fig. 4(a) shows the simulated transmission spectra of the SMC and the DMC. The simulation grid size is 52.3 nm and the resulting spectral resolution is 0.002 nm. Figs. 4(b) and (c) show the simulated WGM intensity distributions of the SMC and the DMC. We can see from Fig. 4(a) that due to the Vernier effect in the DMC where the resonances of each microsphere coincide, the WGM oscillation is strengthened and results in a WGM spectrum with an increased Q factor and a narrowed BW, which is narrowed from 0.495 nm to 0.351 nm, compressed by about 29.3%, as compared to that of the SMC. The ER of the DMC's WGM spectrum is about 45% higher than that of the SMC. Thus, the DMC has a better mode-selection performance than the SMC. Fig. 4(b) shows the WGM intensity distribution of the SMC at 1551.22 nm. Fig. 4(c) shows the WGM intensity distribution of the DMC at 1550.96 nm.

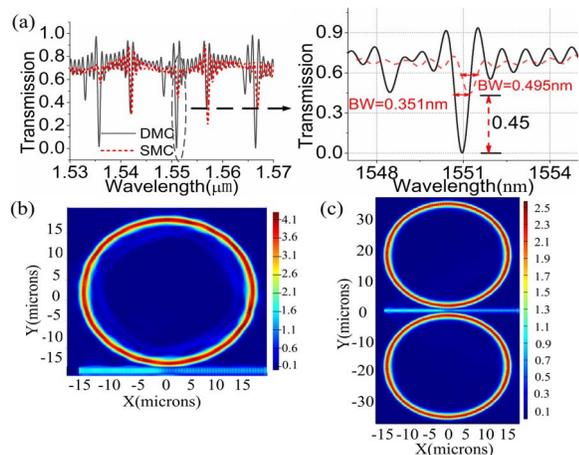


Fig. 4. (a) Comparison of the simulated WGM spectra of the DMC and the SMC. Inset: zoom-in-view WGM spectra. (b), (c) Simulated WGM intensity distributions of (b) the SMC at the wavelength of 1551.22 nm and of (c) the DMC at the wavelength of 1550.96 nm.

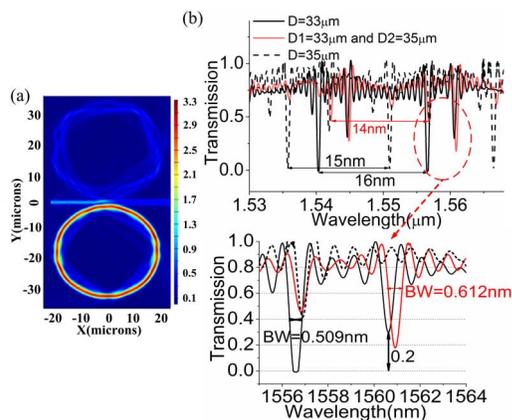


Fig. 5. (a) Simulated WGM intensity distribution of the DMC with different diameters at 1556.81 nm. (b) Comparison of the simulated WGM spectra of the DMC with the same and different D 's. Inset: zoom-in-view WGM spectra.

D. Microsphere Size Dependence for the DMC Resonances

Figure 5(a) shows the simulated WGM intensity distribution of the DMC for microspheres with different diameters ($D_1 = 33 \mu\text{m}$, $D_2 = 35 \mu\text{m}$). The intensity color scale is in linear scale. Different resonant wavelengths between the two microspheres make it difficult to achieve phase matching between the DMC and the taper. In Fig. 5 (a), a WGM at 1556.81 nm is only excited in the microsphere located below the taper. The WGM field intensity is mainly distributed in a single microsphere. Fig. 5(b) shows the comparison of WGM spectra of the DMC, with the microspheres having the same diameters ($D = 33 \mu\text{m}$, $D = 35 \mu\text{m}$) and different diameters ($D_1 = 33 \mu\text{m}$, $D_2 = 35 \mu\text{m}$), respectively. It can be seen clearly that when the two microspheres have the same D of $33 \mu\text{m}$, the WGM spectra have a 20% higher ER, a narrower BW and a larger FSR. Also, it can be seen from Fig. 5(b) that, for the cases that the microspheres are having the same diameters, the smaller $D = 33 \mu\text{m}$ case results in a larger FSR as compared to the larger $D = 35 \mu\text{m}$ case, which is in good accordance with Equation (1).

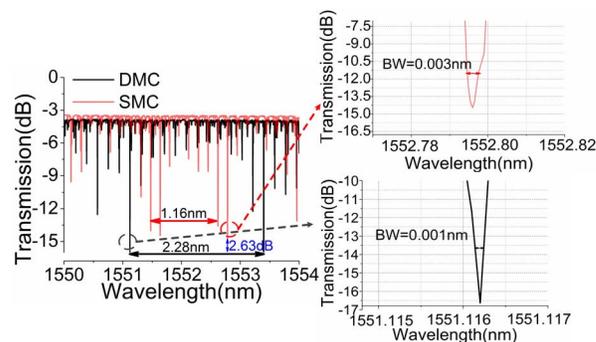


Fig. 6. Comparison of the measured WGM spectra of the SMC and the DMC. Insets: zoom-in-view WGM transmission dips.

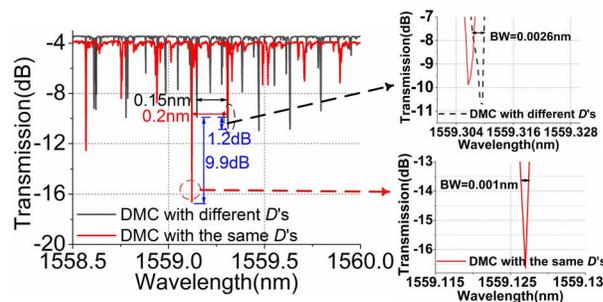


Fig. 7. Measured WGM spectra of the DMC with different and the same diameters. Insets: zoom-in-view WGM transmission dips of the DMC with different D 's (above) and the same D 's (below).

IV. EXPERIMENTAL RESULTS AND DISCUSSIONS

In order to produce stabilized WGM spectra for stable fiber laser operation, over-coupling is applied in our experiments, the microsphere cavities are in touch with the fiber taper with a diameter of about $3 \sim 4 \mu\text{m}$, which is phase-matched well with the microspheres. Figure 6 shows the comparison between the experimentally measured WGM spectra of the SMC and the DMC. The spectral resolution of the transmission measurements is 0.5 pm. The equatorial diameter of the SMC is $\sim 198.4 \mu\text{m}$, and the equatorial diameters of the double microspheres are $\sim 183.4 \mu\text{m}$ and $198.3 \mu\text{m}$, respectively (due to fabrication imperfections, D_1 and D_2 are different by nearly 10%, which is comparable to the fractional difference between $D_1 = 33 \mu\text{m}$ and $D_2 = 35 \mu\text{m}$ in simulations. This may allow the experimental results and the simulation results of the DMC to be qualitatively compared). It can be seen clearly that the DMC can produce a TE-polarized WGM spectrum with a (2.63 dB) higher ER, a narrower BW and a (1.12 nm) larger FSR as compared to those of the SMC, which is also in good accordance with our theoretical analysis.

Figure 7 shows the comparison of the measured WGM spectra within a wavelength span between the DMC with almost the same diameters ($198.4 \mu\text{m}$, $198.3 \mu\text{m}$) and different diameters ($183.4 \mu\text{m}$ and $198.3 \mu\text{m}$). The DMC with the same D 's produces a TE-polarized WGM spectrum with a wider FSR and a higher ER (about 8.7 dB) as compared to the DMC with different D 's (about 1.2 dB) near the wavelength of 1559.31 nm. There are significantly more high-order modes near 1559.1 nm for the DMC with microspheres of

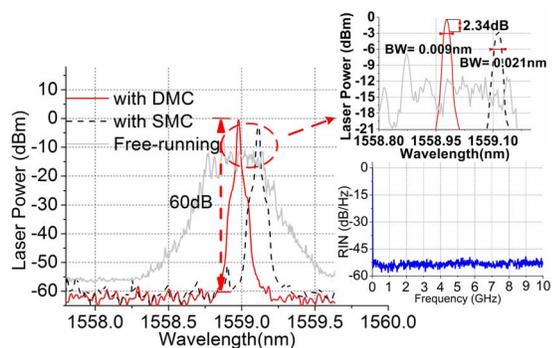


Fig. 8. Comparison of the measured laser spectra with the DMC, the SMC, and the free-running laser operation. Insets: zoom-in-view spectra (above) and measured radio-frequency (RF) spectrum (below) of the proposed single-longitudinal-mode fiber ring laser using self-homodyne method.

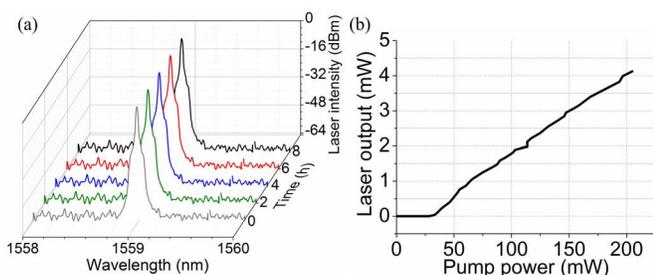


Fig. 9. Measured laser output of: (a) Mode-hopping-free SLM laser output measured for 8 hours with a pump power of 100 mW; (b) Measured laser output versus pump power.

different diameters. Thus, the DMC with the same microsphere diameters has a better mode-selection performance, which is in good accordance with the simulation results.

Figure 8 shows the measured laser spectra based on: the DMC, the SMC and the free-running fiber laser with a pump power of about 100 mW. Using the DMC (with the same parameters as in Fig. (6)) as the mode-selector, we observe the narrowest 3dB-BW of less than 0.01 nm for the SLM laser spectrum, with a highest signal-to-background ratio of about 60 dB. The SLM operation is verified by using a frequency spectrum analyzer (Rohde & Schwarz, FSV30) without beating frequency appeared, as shown in the inset (below). The fiber laser system is placed on a floating vibration-isolated platform and the coupling position is kept stable through a specially designed package. As compared to the SMC method, the DMC method can effectively improve the laser signal-to-background ratio (by 2.34 dB), compress the laser linewidth (by 0.012 nm), and improve the stability (by 8 hr) of the laser output.

Fig. 9 (a) shows the stable SLM laser output measured for 8 hours of using DMC as the mode-selector. When the laser was pumped at about 100 mW, no mode hopping was observed when the laser outputs are monitored and compared at selected times. The wavelength and output power of the

laser is relatively stable. The laser threshold is about 32.2 mW, having a slope efficiency of about 2.4%, as shown in Fig. 9(b).

V. CONCLUSION

We proposed and demonstrated a SLM, narrow-bandwidth fiber laser using the taper-coupled DMC as the mode-selector inside a fiber ring cavity. As compared with the SMC, the DMC produces WGM spectra with a higher ER, a higher SMSR, a narrower BW as well as a larger FSR, which achieves a SLM and a highly stabilized laser output with a narrower BW of < 0.01 nm, a higher signal-to-background ratio of about 60 dB, and a higher laser efficiency as compared to using the SMC for mode selection. This fiber laser could find potential applications in optical communications and sensing systems.

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